



Dynamical effects of radiative losses in high-mass microquasars

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Outline

1 Microquasars

2 Methods

- Code
- Model description

3 Results

- Dynamical effects
- Emissivity

Plan

1 Microquasars

2 Methods

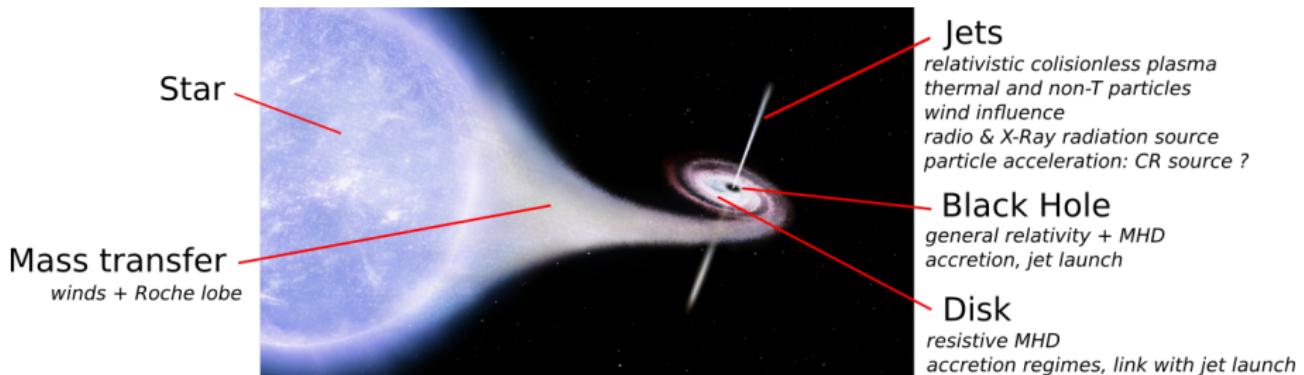
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High Mass Microquasars

Lab of relativistic plasma physics:

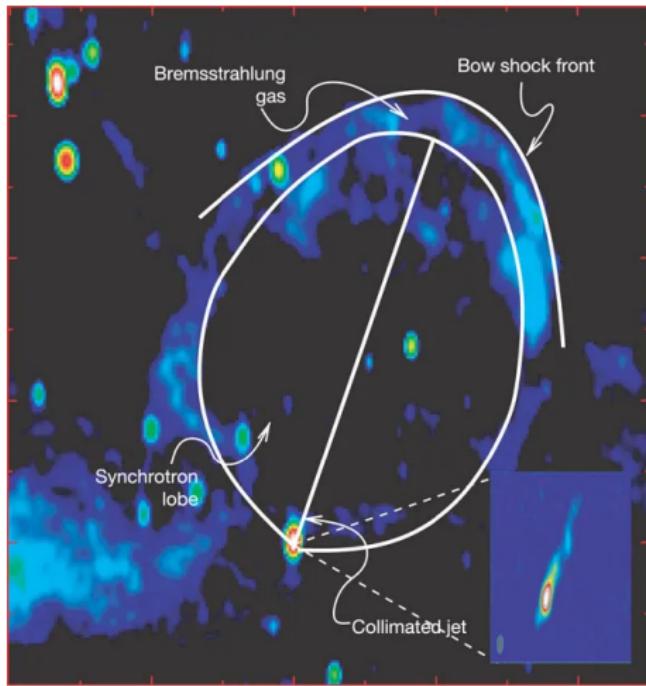


2 main spectral states:

Low Hard State (LHS): peak ~ 100 keV + radio jets (Corbel et al. 2012)

High Soft State (HSS): black body-like peak 1 keV

Observed jets



large-scale structure
Different zones = different radiative processes

Collimated jet misaligned with structure:
are we sure it originates from the jet ?

Cygnus X-1 @1.4 GHz (Gallo *et al.* 2005). Structure spans \sim 15 ly

Motivations

Questions:

- jet contribution to HE spectrum ?
- dynamical effects of wind & radiative losses ?

Difficulties:

- midly relativistic regime: usual approximations not valid
- adding radiative processes = calculation time++

State of the art:

- no global simulations with dynamical losses

My work:

- calculation of radiative processes in our general regime
- RHD and radiative losses in A-MaZe
- global simulations over a long time period

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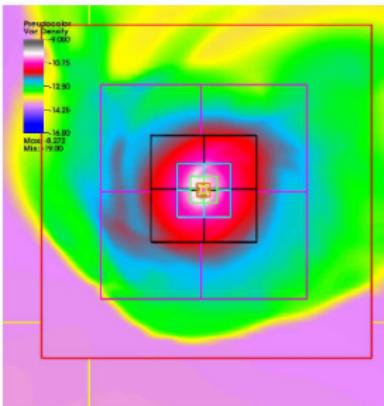
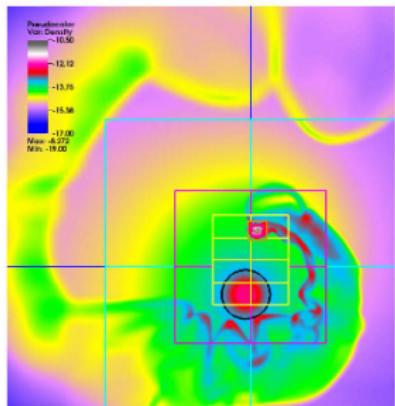
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Code

A_MaZe (R. Walder, D. Folini)



newtonian hydrodynamics
conservative scheme
finite volumes
multiD *upwind* algorithm
AMR

accretion in Cygnus X-1 (Walder et al. 2014)

Box dimension 1AU (left), ~ 0.01 AU (right), refinement up to $8R_s$ (right)

Relativistic hydrodynamics

Primitive variables:

rest frame density	ρ
velocity	v^i
pressure	p

Conservative variables:

$$\begin{aligned} D &= \gamma\rho \\ S^i &= cT^{0i} = \rho h\gamma^2 v^i \\ \tau &= T^{00} - J^0 c^2 = \rho h\gamma^2 - p \end{aligned}$$

Conservation of current $J^\mu = \rho u^\mu$ and energy-momentum tensor
 $T^{\mu\nu} = \rho h u^\mu u^\nu + pg^{\mu\nu}$

$$\begin{aligned} \nabla_\mu J^\mu &= 0 \\ \nabla_\mu T^{\mu\nu} &= 0 \end{aligned} \quad \Rightarrow$$

$$\partial_t D + \partial_i(Dv^i) = 0 \quad (1)$$

$$\partial_t S^j + \partial_i(S^j v^i + pc^2 \delta^{ij}) = 0 \quad (2)$$

$$\partial_t \tau + \partial_i S^i = -P_{loss} \quad (3)$$

Radiative mechanism

$$P_{loss} = P_{Bremsstrahlung} + P_{Synchrotron} + P_{InverseCompton}$$

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Bremsstrahlung + **relativistic correction** (Rybicki & Lightman 1979):

$$P_{Bremsstrahlung} = 1.4 \cdot 10^{-27} T^{1/2} n_e n_p \bar{g}_B (1 + 4.4 \cdot 10^{-10} T) \quad (4)$$

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Synchrotron and IC: RL79 modified with **relativistic Maxwellian distribution**:

$$P = \frac{4}{3} \sigma_T c n_e \gamma^2 \beta^2 U \quad \Rightarrow \quad P = \frac{4}{3} \sigma_T c n_e \gamma \Theta \frac{K_3(1/\Theta)}{K_2(1/\Theta)} U$$

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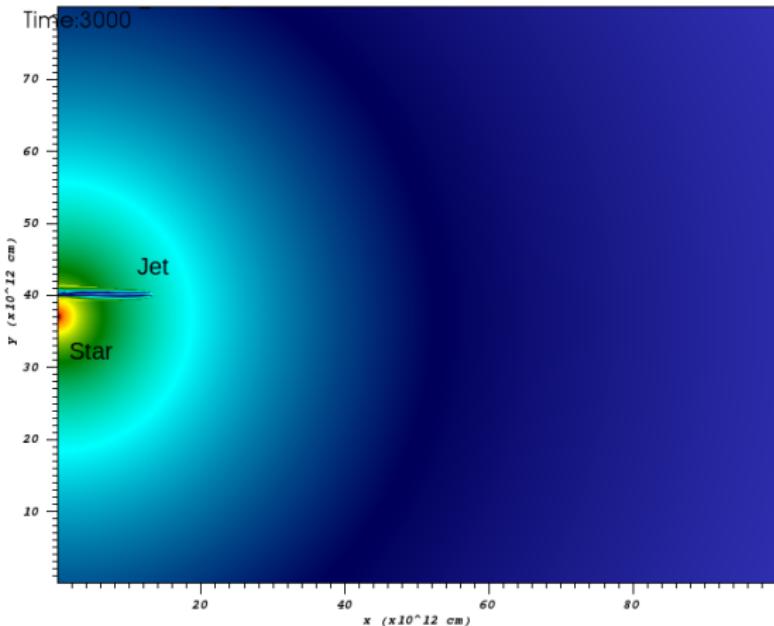
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Large scale



full-size density map
Star - BH separation $a = 0.2$ AU

reflexive conditions at $x=0$, outflow elsewhere

ρ_j, v_j, T_j at injection

box $\sim 7 \times 5 \times 5$ AU

coarse grid: $250 \times 200 \times 200$,
 $dx = dy = dz = 4 \cdot 10^9$ m

5 levels up to $\times 64$ refinement
fixed grid centered on jet for performance

12 setups

following figures all in log scale

Jet morphology

Cocoon (1): shocked wind material, **large emitting bubble**

Inner cocoon (2): jet + wind material, **instabilities**

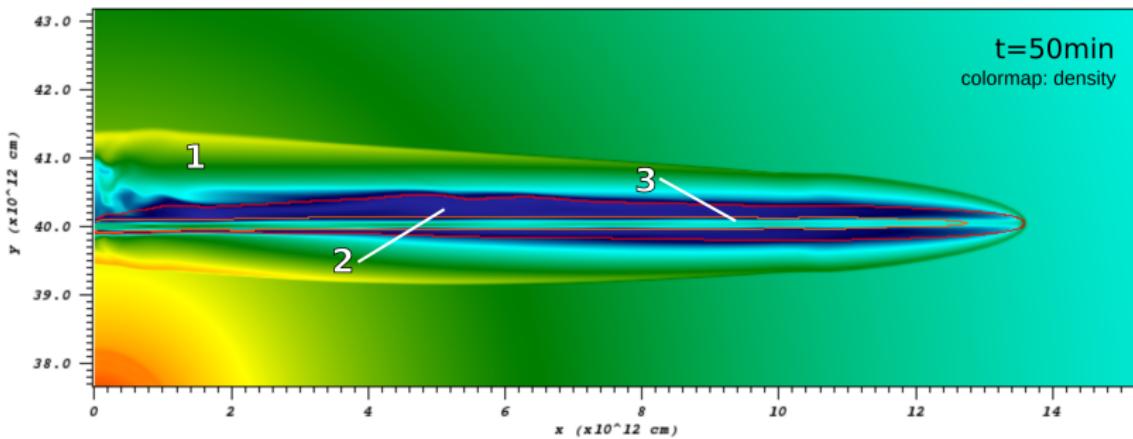
Beam (3): injected matter, **recollimation shocks**

Here:

$$10^{-17} < \rho < 10^{-13} \text{ g cm}^{-3}$$

$$10^7 < T < 10^{11} \text{ K}$$

P_j	ρ_j	v_j	r_0		\dot{M}_*	v_w	T_*	(cgs)
10^{37}	$1.3 \cdot 10^{-15}$	$1 \cdot 10^{10}$	$5 \cdot 10^{10}$		$3 \cdot 10^{-6}$	10^8	$3 \cdot 10^4$	



Jet morphology

Cocoon (1): shocked wind material, **large emitting bubble**

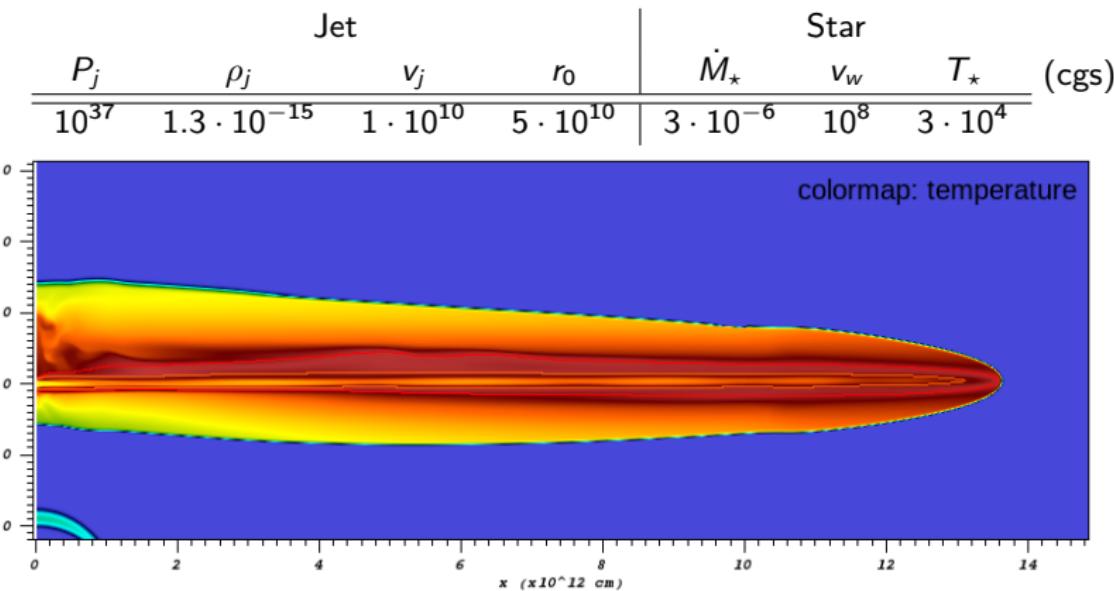
Inner cocoon (2): jet + wind material, **instabilities**

Beam (3): injected matter, **recollimation shocks**

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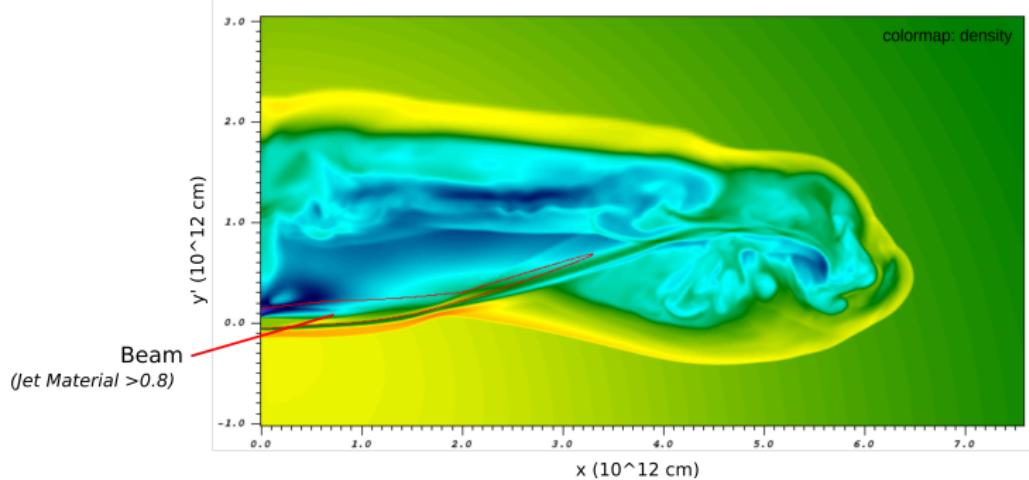
$$10^{-17} < \rho < 10^{-13} \text{ g cm}^{-3}$$

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Jet bending

$$\frac{P_j}{10^{36}} \quad \rho_j \quad v_j \quad | \quad v_w \quad (\text{cgs}), \quad t=1\text{h}40$$



$P_j < P_{lim}$, strong bending of the beam & instabilities (Molina & Bosch-Ramon 2018)
 → misaligned large-scale structure of Cyg X-1 ?

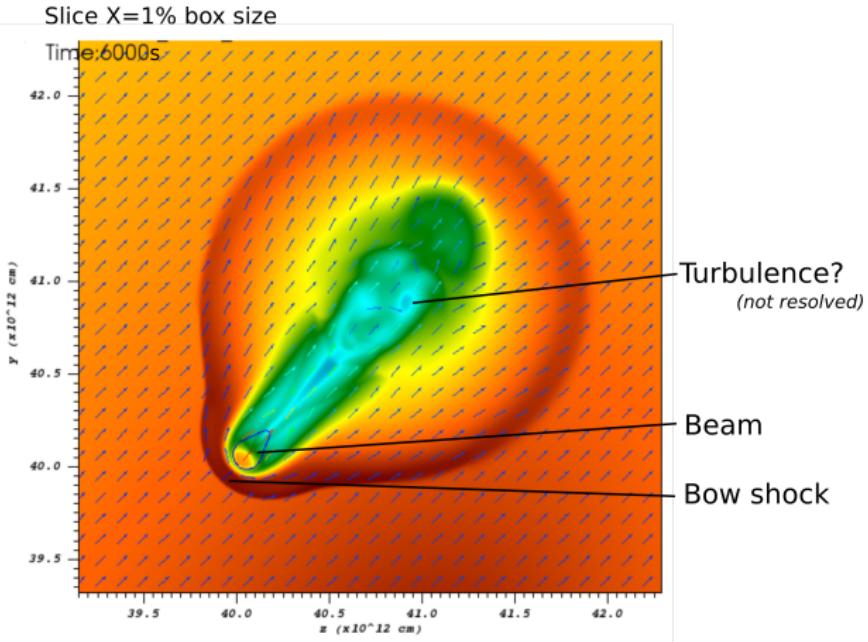
Cocoon deformation

$$\frac{P_j}{10^{36}} \quad \rho_j \quad v_j \quad | \quad v_w \quad (cgs), t=1h40$$

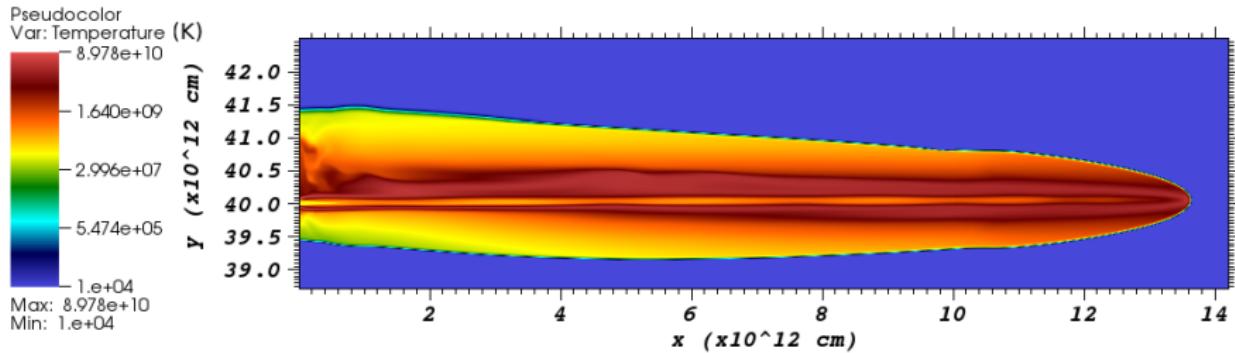
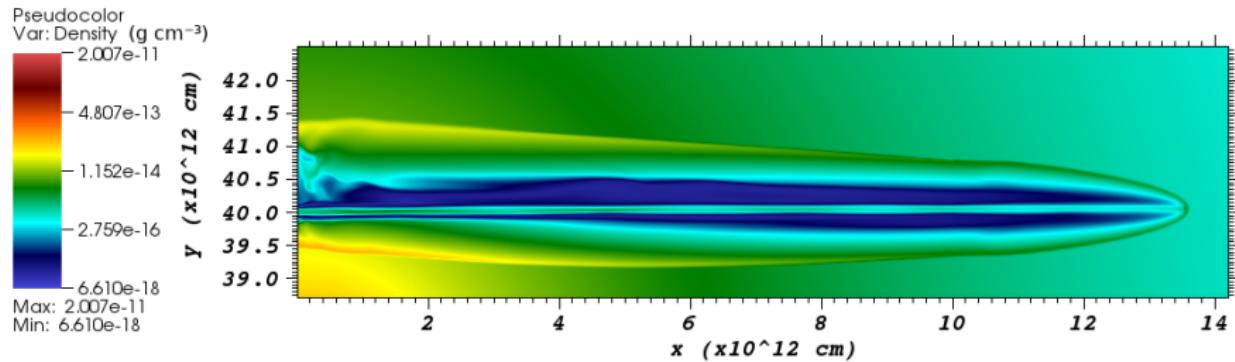
$$4.72 \cdot 10^{-15} \quad 3 \cdot 10^9 \quad | \quad 1.5 \cdot 10^8$$

Previous simulation sliced
at $x = a/3$:
strong stellar wind

Bow shock around beam,
cocoon pushed by wind

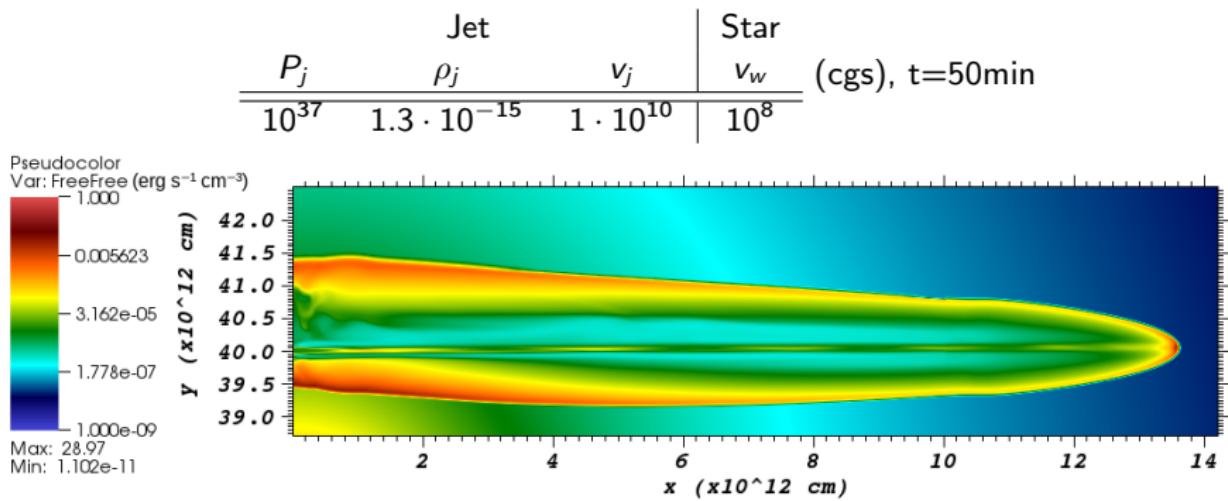


Maps



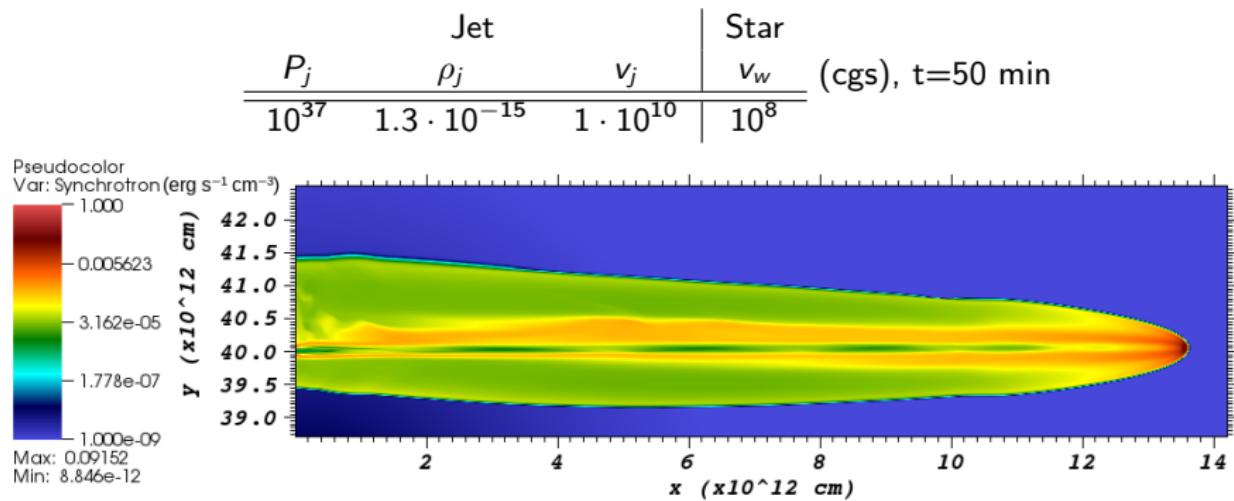
logscale maps of ρ and T

Bremsstrahlung



$\epsilon \propto \rho^2 T$, strong at front shock and cocoon interface

Synchrotron

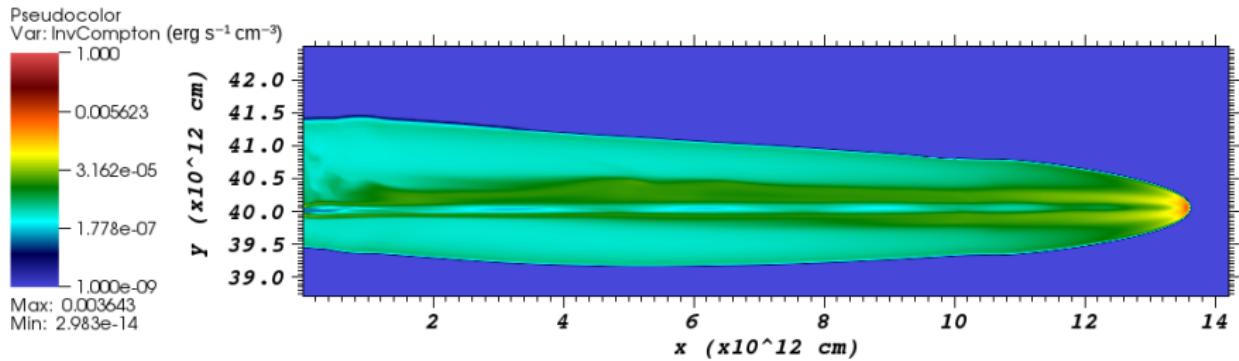


Strongest at shock front

$$U = B^2/8\pi, \text{ with } B_{*,0} = B_{j,0} = 10 \text{ G}$$

Inverse Compton

Jet			Star	(cgs), t=50 min
P_j	ρ_j	v_j	v_w	
10^{37}	$1.3 \cdot 10^{-15}$	$1 \cdot 10^{10}$	10^8	

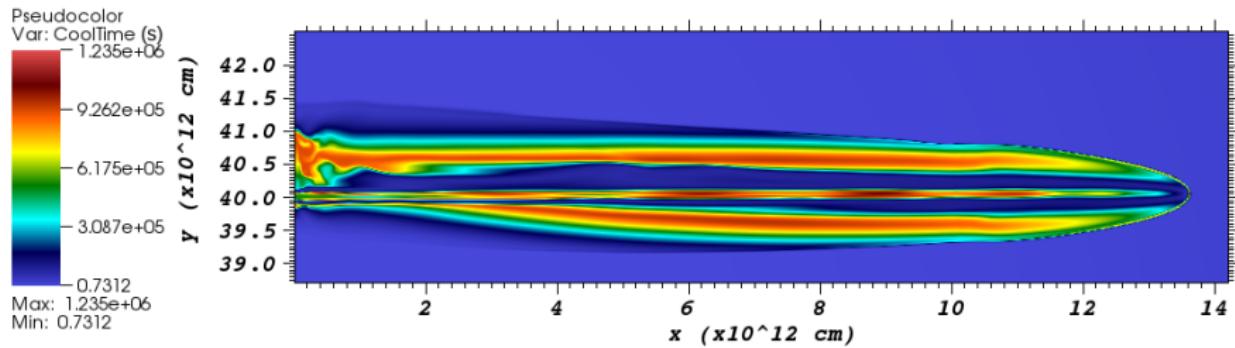


Same ρ , T dependence as synchrotron \rightarrow same structure, overall weaker
 $U \propto \sigma_T T_*^4$

Cooling time

$$t_{cool} = e_{int}/\sum \epsilon_i$$

Jet			Star	(cgs), t=50 min
P_j	ρ_j	v_j	v_w	
10^{37}	$1.3 \cdot 10^{-15}$	$1 \cdot 10^{10}$	10^8	



low t_{cool} near beam injection (stronger winds) and inner cocoon
high t_{cool} for beam

Conclusions and prospects

Conclusions

implementation of SR hydro with radiative losses in A_MaZe

jet bending under critical P_j and cocoon deformation

high t_{cool} variability: modified structure but no effects on propagation (yet)

Prospects

refine the model:

$T > 10^{10}$ K (e^\pm production), adapt the EoS

add non-thermal $e^- \rightarrow$ modifications to the dynamics ?

launch a simulation at best precision

draw radiation maps (radio/UV/X) and spectras

simulate other systems